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# A DFT Study of Structural and Electronic Properties of ZnS Polymorphs and its Pressure-Induced Phase Transitions

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A systematic first-principles investigation, by using the density functional formalism with the non-local B3LYP approximation including a long-range dispersion correction, has been performed to calculate the structural and electronic properties and phase transitions under pressure of the three phases of ZnS (cubic zinc blende, ZB, hexagonal wurtzite, W, and cubic rock salt, RS). Numerical and analytical fittings have been carried out to determine the equilibrium unit cell geometry and equation of state parameters for the ZnS phases. The band structures, energy gap, density of states, and vibrational frequencies and their pressure dependences are investigated. The present results illustrate that both phases, W and ZB, present very similar enthalpy and the RS phase becomes thermodynamically more stable than ZB and W structures at 15.0 and 15.5 GPa, respectively. These phase transitions are accompanied by an increase of the first shell coordination number of Zn atom and by a cell volume collapse of 13.9 and 14.3% for ZB and W phases, respectively. The atomic contributions of the conduction and valence bands, as well the binding energy for the Zn 3d orbital have been obtained.

## I. INTRODUCTION

Zinc sulfide (ZnS) has attracted much attention as an important semiconductor material because of its wide range of applications in photonics and its strength for novel applications, including light-emitting diodes (LEDs), electroluminescence, lasers, infrared windows, flat panel displays, sensors, catalysis, etc.<sup>1-5</sup>

Understanding the relationship between structural and electronic properties of the semiconductor materials and their pressure-induced phase transitions is crucial to design novel nanostructures with tunable properties. ZnS presents three polymorphs: cubic zinc blende (ZB), hexagonal wurtzite (W) or

the rarely observed cubic rock salt (RS).<sup>6</sup> Each phase has unique physical properties, for instance, different lattice vibration properties and nonlinear optical coefficients.<sup>7-9</sup> In particular, ZB is the thermodynamically most stable ZnS phase, while the W polymorph is stable above 1293 K.<sup>9</sup> However, RS structure can only be obtained at relatively high pressures.<sup>10,11</sup> ZB and W phases have industrial applications, and due their size- and shape-dependent properties both materials are capable of being obtained in different ways in order to tune their properties to specific needs.<sup>9,12,13</sup> In this context, with decreasing the particle size, the relative stability of the two phases changes and a low temperature synthesis of ZnS nanoparticles with W structure has been reported.<sup>6,13-15</sup>

Several routes are described in the literature to obtain ZnS with different morphologies<sup>6,16-23</sup> and a considerable effort has been devoted to characterized these materials by means of different techniques such as x-ray diffraction (XRD),<sup>17,18</sup> extended X-ray absorption fine structure,<sup>19</sup> high-resolution transmission electron microscopy (HRTEM),<sup>20</sup> Fourier transform infrared (FT-IR) and Raman spectroscopy,<sup>8-21</sup> X-ray photoelectron spectroscopy (XPS)<sup>6</sup> and photoluminescence (PL) measurements.<sup>6,22,23</sup>

Theoretical and computational studies, based on the Density Functional Theory (DFT)<sup>24,25</sup> have emerged to provide important information regarding the electronic and structural properties of solid materials and molecules,<sup>24-29</sup> and have shown great value, not only in the interpretation of experiments, but also in the prediction of important aspects of new properties and in the design of new devices. Nevertheless, despite the recent improvements in DFT, there are still difficulties in using DFT to properly describe van der Waals interactions (dispersion) and charge transfer excitations due to the lack of exact Hartree–Fock exchange in some functionals.<sup>30-34</sup> To overcome these limitations, an empirical correction to include dispersion interactions in DFT methods has been proposed by Grimme.<sup>35-37</sup> It is a damped pairwise London-type term of the form  $C_6R^{-6}$  and this correction is applied to calculate

more accuracy values of energies and gradients in the simulation of periodic models, and it is implemented in CRYSTAL 09 code.<sup>38</sup>

Among the most interesting phenomena, the pressure-induced ZnS polymorphism (ZB, W and RS) is specially relevant and has been a subject of many theoretical and experimental studies to understand the observed changes under pressure.<sup>11,17,39-41</sup> However, understanding this polymorphism at atomic level is very important for technological applications of materials based on ZnS and computational simulations can shed some light on to this important issue. Knowledge of the pressure-dependent phase stability and the relationship between physical/chemical properties and crystal/electronic structures can offer the way to systematically search new complex metal oxides. In this sense, this work can be considered as a prolongation of previous high pressure studies of pressure induced structural phase transitions of different binary and ternary metal oxides.<sup>42,43</sup>

In this work, we report periodic first-principle calculations based on DFT to obtain the energy, geometry, and vibrational frequency by including the empirical dispersion correction on ZnS polymorphs and its pressure-induced phase transitions. The paper is organized as follows. Section 2 describes the computational details. In section 3, we present our theoretical results together with the discussion concerning the local compressibility and phase stability, as well as the electronic structure analysis for the different phases; in addition, vibrational properties as derived from calculations are discussed in detail. Finally, we summarize our main conclusions in Section 4.

## II. COMPUTATIONAL METHODS

In the present study, we use periodic DFT calculations with the B3LYP-D hybrid functional<sup>44,45</sup> with a long-range dispersion correction, as implemented in the CRYSTAL09 computer code.<sup>38</sup> CRYSTAL 09 uses Gaussian-type basis set to represent crystalline orbitals as a linear combination of Bloch functions defined in terms of local functions (atomic orbitals).<sup>46</sup> A long-range dispersion correction proposed by

Grimme has been considered in the calculations.<sup>35-37</sup> The zinc and sulfur atomic centers were described by all-electron 6-31G\* and 86-311 G\* basis sets, respectively. Basis sets were optimized in this study and are available in supplementary material. Optimization of the exponents for the outermost *s* and *d* shells was carried out to minimize the total energy of the structure at experimental parameters ( $\alpha_{sp(Zn)} = 0.14349998$ ,  $\alpha_{d(Zn)} = 0.73000001$  and  $\alpha_{sp(S)} = 0.38000002$ ).<sup>22</sup> The Powell algorithm method<sup>47</sup> was used to perform the optimization procedure of the basis sets. The level of calculation accuracy for the Coulomb and exchange series was controlled by five thresholds set to ( $10^{-8}$ ,  $10^{-8}$ ,  $10^{-8}$ ,  $10^{-8}$ , and  $10^{-18}$ ). The shrinking (Monkhorst–Pack)<sup>48</sup> factor was set to 6, which corresponds to 80 independent *k*-points in the irreducible part of the Brillouin zone integration.

Conventional unit cells of the three polymorphs belong to the *F-43m* (ZB), *P6<sub>3</sub>mc* (W) and *Fm-3m* (RS) spaces groups (see Figure 1). Lattice parameters, for ZB and RS structures, and also atomic positions for the W structure were fully optimized. The lattice parameters and atomic positions obtained by means by optimized structure in zero pressure equilibrium properties for ZnS polymorphs were used in the Diamond Crystal and Molecular Structure Visualization program (Version 3.2f for Windows)<sup>49</sup> to carry out the XRD simulation. To take into account the effect of pressure on this system, all the geometric parameters (*a*, *c* and *u*) were optimized for ZnS polymorphs. Fittings with a Birch–Murnaghan third-order equation of state (EOS) of the computed energy–volume data provide values of the zero-pressure bulk modulus (*B*<sub>0</sub>) and its pressure derivative (*B*<sub>0</sub>') as well as enthalpy–pressure curves for the polymorphs studied.<sup>50</sup>

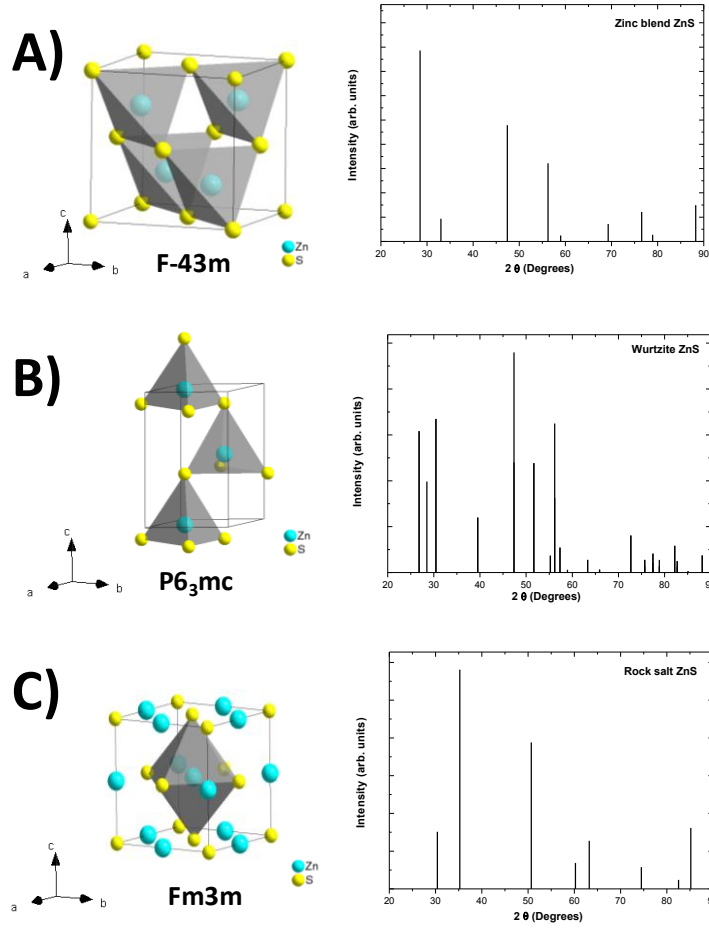
The characteristics and properties of ZnS polymorphs (ZB, W and RS) was evaluated and discussed in terms of charges and deformations induced by pressure on the ZnS clusters. The XcrysDen program<sup>51</sup> was used for the band structure drawing design. The analysis of the vibrational modes and their corresponding frequencies were calculated through numerical second derivatives of the total energies as implemented in the CRYSTAL09 package.

### III. RESULTS AND DISCUSSION

The structure of the three phases of ZnS and their respective XRD simulation are shown in Figure 1. These results are in very good agreement with other studies<sup>5,6,22,52-55</sup> The theoretical lattice parameters and other theoretical and experimental data are displayed in Table 1. In W and ZB structures each Zn atom is surrounded by four S atoms forming tetrahedral clusters  $[\text{ZnS}_4]$  which share one or two corners, respectively. W phase is more compact than ZB phase, and the W phase shows a lack of a center of symmetry which generates a residual polarization in this  $[\text{ZnS}_4]$  cluster. The optimized W structure belongs to the hexagonal space group  $P6_3mc$  with the lattice parameters:  $a = 3.79$  and  $c = 6.14$  Å and contains two formula units of ZnS per unit cell. The ratio of  $c/a = 1.621$  deviates slightly from the value of  $c/a = \sqrt{8/3} = 1.633$  for the hexagonal phase of ZnS.<sup>56</sup> For RS phase in the ZnS crystal, each Zn atom is surrounded by six S atoms at the corners of the  $[\text{ZnS}_6]$  octahedral clusters. The structural characteristics of these clusters, as building blocks of this system, may explain some particular properties of ZnS crystal.

**Table 1:** Comparative results between the structural parameters and unit cell volume of ZnS polymorphs in the equilibrium configuration obtained in this work with those published in the literature.

Polymorphs	a(Å)	c(Å)	u	Ref.
ZB	5.42	-	0.250	This work
	5.33	-	-	39
	5.58	-	-	41
	5.42	-	-	53
W	3.79	6.14	0.379	This work
	3.82	6.26	-	4
	3.82	6.26	0.375	52
	3.83	6.26	0.379	22
RS	5.09	-	0.500	This work
	5.01	-	-	39
	5.21	-	-	41
	5.02	-	-	54



**Figure 1:** Schematic representation of the unit cells (left) and simulated XRD patterns (right) for the ZnS phases: a) Zinc blende, b) Wurtzite and c) Rock salt.

Raman spectroscopy is a well known and useful method for investigating the behavior of symmetry changes in semiconductor materials.<sup>21,57-60</sup> The theoretical Raman-active modes are shown in the Table 2. ZB structure is a polar crystal without an inversion center, and it belongs to the point group F-43m with only four particles per unit cell and has three optical modes triply degenerate with symmetry ( $\Gamma_{15}(F)$ ) and both of which are infrared and Raman active.<sup>21,59</sup> In particular, the experimental Raman spectrum for ZB structure has been reported by Nilsen,<sup>59</sup> and according to these results, two bands appear that are assigned the transverse optic (TO) mode at 271  $\text{cm}^{-1}$  and the longitudinal optic (LO) mode with a higher frequency at 352  $\text{cm}^{-1}$ . The symmetry of these modes and the predicted frequencies

for ZnS polymorphs is well established by their polarization characteristics.<sup>8,21,59</sup> The W structure belongs to the point group  $P6_3mc$ , with only two particle per unit cell and has six optical modes ( $\Gamma = A_1 + 2B_1 + E_1 + 2E_2$ ), being the  $A_1$ ,  $E_1$  and  $E_2$  Raman active and  $A_1$  and  $E_1$  modes infrared active.<sup>8,21,59</sup>  $B_1$  modes are silent modes in both types of spectra.<sup>8</sup>

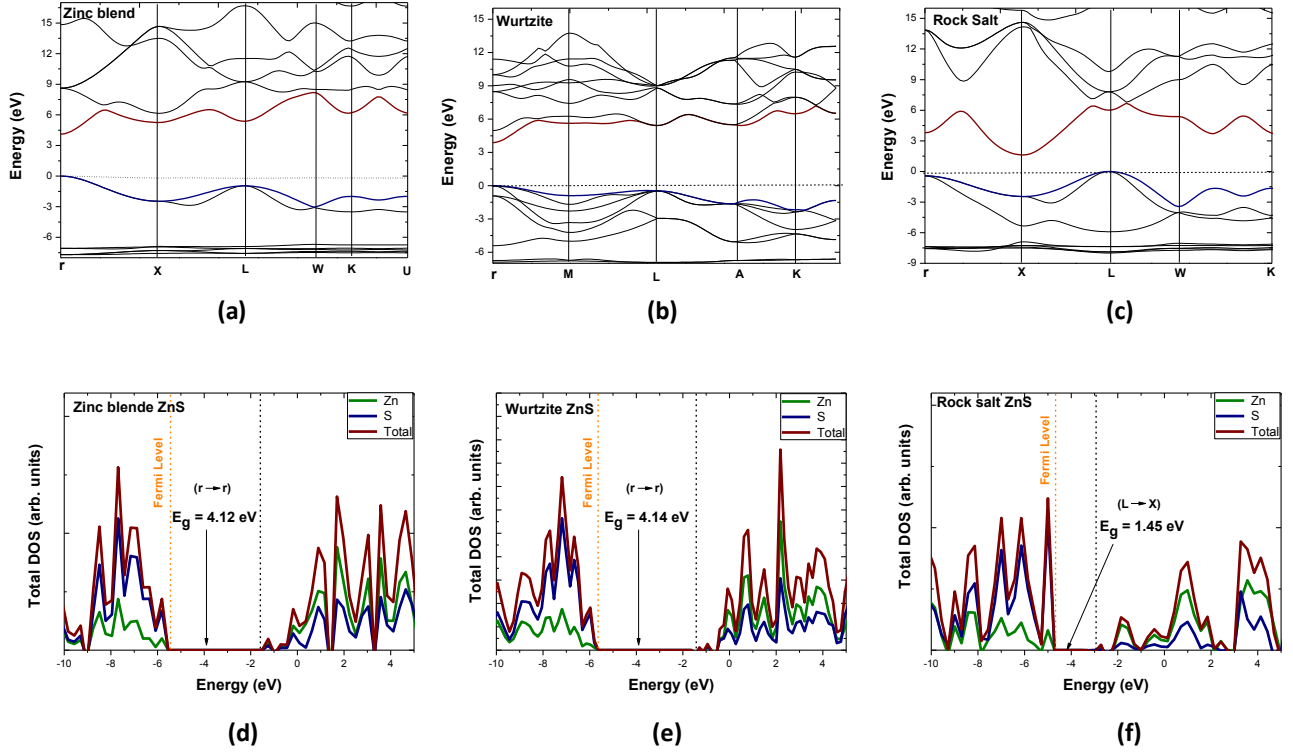
**Table 2:** Comparative results between the experimental and theoretical Raman modes of ZnS polymorphs in  $\text{cm}^{-1}$ .

ZB structure	This work	Ref. 21	Ref. 58
F	328.01	277 (TO) 340 (LO)	271 (TO) 352 (LO)
W structure	This work	Ref. 21	Ref. 8
$E_2$	81.70	76	72
A	315.06	287	273
$E_1$	328.34	288	273
$E_2$	333.08	296	286

The band structure and DOS of these materials were calculated and the results are shown in Figure 2. An analysis of the band structure (see Figure 2 (a-c)), shows that both ZB and W phase have a direct band gap at the  $\Gamma$  point, while RS phase has an indirect band gap at the  $L \rightarrow X$  point. The calculated band gaps for ZnS polymorphs are 4.10, 4.14 and 1.45 eV for ZB, W and RS respectively, in good accordance with the values reported in the literature.<sup>6,22,39</sup> Theoretical calculations for the W phase were reported previously,<sup>22</sup> not including long-range corrections using the same theoretical level. The reported theoretical value of the band gap for the W phase was 3.88 eV while in present study, there is an increase in the estimated value for the band gap of 4.14 eV. These results show that the inclusion of dispersion in the DFT calculation generates an increase of approximately 6.3% in band gap value for the phase W. Similar trends were observed for other phases.



An analysis of the DOS for the bulk ZnS polymorphs shows that uppermost valence band (VB) consists mainly of S  $3p$  orbitals with a lesser contribution of hybrids Zn  $4s$  and  $4p$  orbital states. The presence of Zn  $3d$  states in the VB in the ZnS polymorphs models reveals a strong bonding character between S and Zn. The conduction band (CB) consists mainly of hybridization of Zn  $4s$  and  $4p$  orbital states with a small contribution of S  $3p$  orbitals. This results is also found in the DOS analysis of Zn-based II–VI semiconductor materials.<sup>61</sup> In addition, we have calculated the binding energies (BEs), as the difference  $\Delta E = E^{\text{HOMO}} - E^{\text{core-level}}$ , of the Zn  $3d$  orbitals for the three ZnS polymorphs, and the results are as follow: ZB, W and RS structures of 7.43, 6.98 and 7.67 eV, respectively.

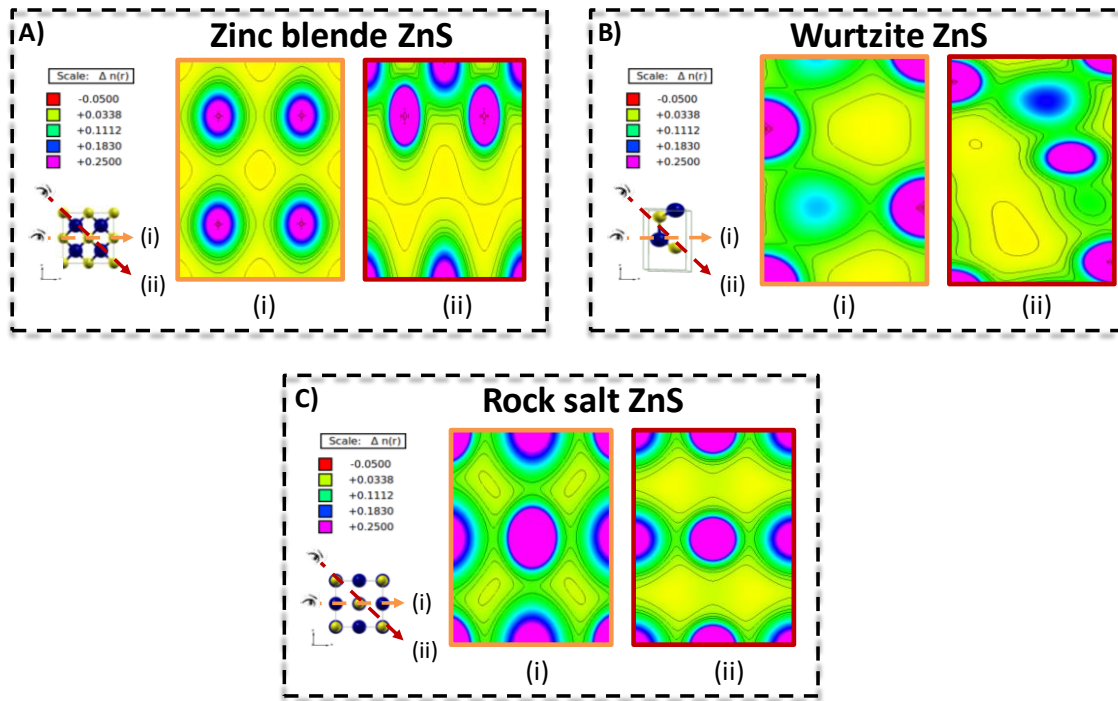


**Figure 2:** Band structures and Projected DOS on atomic levels for ZnS polymorphs: (a-d) Zinc blende, (b-e) Wurtzite and (c-f) Rock salt.

The atomic contributions of the VB and CB in DOS analysis, named as atoms present at the active site (APAS), have been obtained following a recently developed protocol; for more details on this methodology see reference,<sup>62,63</sup> and the results for ZnS polymorphs are presented in Table 3. Our results can reveal the fundamental relationship between the nature of chemical bonding and bulk properties in ZnS polymorphs, showing the presence of mixed ionic-covalent bonding for these materials. The bonding is predominantly covalent for ZB and W phases while it is more ionic for RS structure. The Zn–S distance is 2.308 Å, 2.313/2.316 Å and 2.510 Å for ZB, W and RS, respectively. Figure 3, the charge density contours are depicted for bulk ZnS polymorphs. A detailed analysis of these results point out that a larger contribution of covalent bond in the ZnS clusters causes an energetic stabilization of the crystal structure. These results are in good agreement with previous studies reported by Jaffe et al.<sup>41</sup>

**Table 3:** APAS contribution (%) for the density of states in ZnS polymorphs.

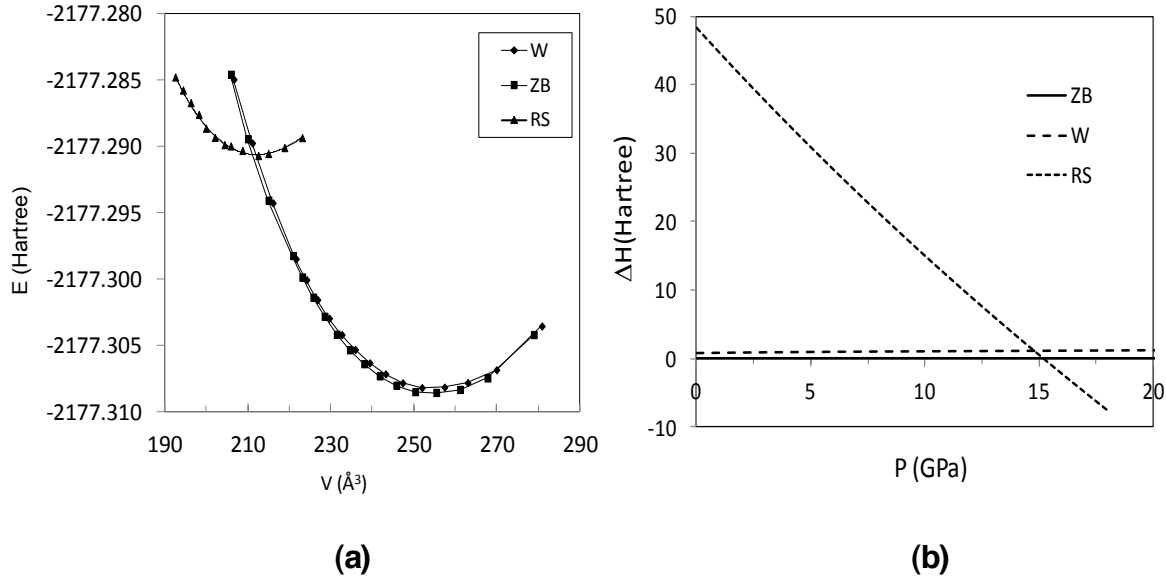
Polymorphs	atoms	APAS contribution %	
		VB	CB
ZB	Zn	30.16	63.05
	S	69.84	36.95
W	Zn	29.90	61.65
	S	70.10	38.35
RS	Zn	31.90	71.58
	S	68.10	28.42



**Figure 3:** Charge density maps in ZnS polymorphs: a) Zinc blende, b) Wurtzite and c) Rock salt.

Calculated lattice parameters and distances between the Zn and S atoms at different pressures as well as the bulk modulus and its pressure derivative for the three phases of ZnS structure are listed in Table 4. For the ZB phase we have obtained the EOS being  $B_0 = 115.6$  GPa and  $B_0' = 3.3$ , higher than the experimental value reported by Nazzari et al.,<sup>54</sup> of 83.3 GPa. For the W phase, the bulk modulus was  $B_0 = 112.3$  GPa and  $B_0' = 3.5$ , a larger value than other experimental<sup>64</sup> and theoretical studies.<sup>39</sup> Regarding the polyhedral changes taking place as the pressure is applied, calculations show that the distortion of the  $[\text{ZnS}_4]$  tetrahedral clusters is reduced in the W phase, and at 35 GPa there is only one distance Zn-S.

Figure 4(a) shows the energy versus volume curves for the three different structures of ZnS. The enthalpy curves are plotted as a function of pressure in Figure 4(b). Both phases, W and ZB present very similar enthalpy. Upon further compression we found that rock-salt structure become thermodynamically more stable than ZB and W at 15 and 15.5 GPa, respectively. These transition pressures are in excellent agreement with the experimental values of 15.0-16.0 GPa.<sup>65,66</sup>



**Figure 4:** (a) Energy versus volume curves for the three different structures of ZnS. (b) The enthalpy curves as a function of pressure.

The transition to RS involves a larger volume collapse. For the RS phase the  $B_0 = 155.4$  GPa and  $B_0' = 5.1$ , higher than the experimental value reported by Nazzari et al.<sup>54</sup> of 104.4 GPa, a bulk modulus approximately 25% larger than the value for ZB and W structures. This decrease of the bulk compressibility is caused by the rearrangement of the polyhedral units that takes place at the transition.

The W-to-RS phase transition is a reconstructive phase transition and the transformation mechanism is associated with large atomic displacements and strain. The volume reduction during the phase transition for both ZB and W phases is predicted to be about 13.9 and 14.3%, respectively, which are comparable with the experimental result of 15.7–17%,<sup>64</sup> and the theoretical result of 14.4%.<sup>67</sup> Despite the controversy about the value of transition pressure obtained by experiments and theory, both clearly demonstrate that the transition pressures of the W-to-RS and ZB-to-RS phase changes are very close to each other. From enthalpy calculations, we can state that no phase transition between the bulk ZB and W phases occurs with the application of pressure alone because their enthalpies are always parallel to each

other. However, this is not the case for W nanocrystals as clearly demonstrated in experiment.<sup>17,68</sup> This behavior of nanocrystals is probably associated with the presence of significant surface energy effects.

The average Zn-S bond distances increased in 8.8% and 9.0% at the transition pressure W-to-RS and ZB-to-RS of 15.0 and 15.5 GPa, respectively. Table 4 shows the pressure evolution of the unit-cell parameters of ZnS phases. An analysis of Table 4 shows that the variation with pressure of the  $c$  parameter is higher than for the  $a$  parameter for the W structure. However, the computed linear compressibilities of  $1.3 \times 10^{-3} \text{ GPa}^{-1}$  and  $1.7 \times 10^{-3} \text{ GPa}^{-1}$  for  $\kappa_c$  and  $\kappa_a$ , respectively, show a slightly anisotropic behaviour for W structure. From these numbers we recover  $B_0$  values consistent with those deduced from  $p$ - $V$  data. Therefore, W structure is not significantly compressed along  $c$ -direction, according to previous work of Durandurdu et al.,<sup>39</sup> in which it is reported a W-to-RS mechanism different from other hexagonal and tetragonal paths observed in a variety of materials with W structure. The ZB-to-RS transformation mechanism of ZnS is well established in the literature<sup>57,69-71</sup> by means of an orthorhombic pathway (via an intermediate state with  $Pmm2$  symmetry). In particular, the shape of the cluster is determined by the contribution of covalent binding, preferring the tetrahedral coordination in the ZnS clusters for ZB and W structures. On the other hand, with increasing of the Zn-S bond ionic character, as is evidenced for RS structure, an octahedral coordination of the ZnS clusters is favored.

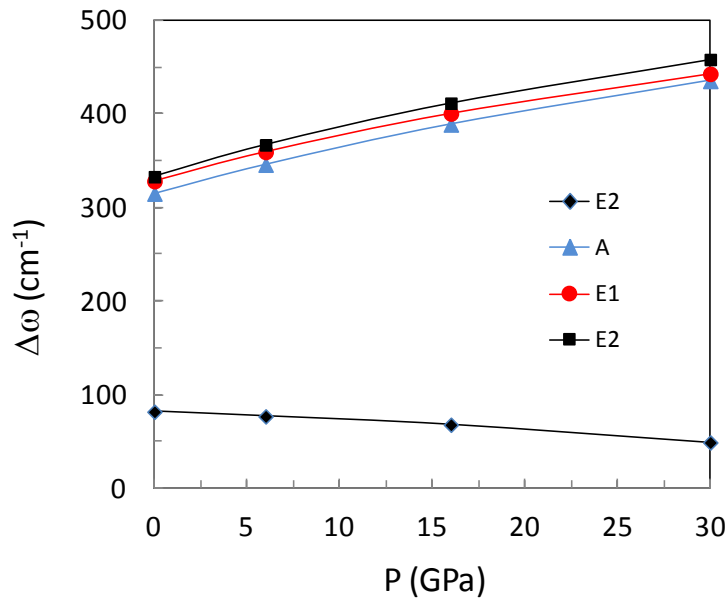
**Table 4:** Calculated lattice parameters from primitive unit cells and distances between the Zn and S atoms at different pressures as well as the bulk modulus and its pressure derivative for the three phases of ZnS structure.

Zinc Blende ( $B_0 = 115.6$ GPa and $B_0' = 3.3$ )				
P (GPa)	a (Å)		V(Å <sup>3</sup> )	Zn-S (Å)
0.0	3.77		37.84	2.308
4.0	3.72		36.42	2.278
8.0	3.68		35.27	2.254
12.0	3.65		34.29	2.233
16.0	3.62		33.45	2.215
20.0	3.59		32.73	2.199
25.0	3.56		31.85	2.179
30.0	3.53		31.12	2.162
35.0	3.51		30.51	2.148
Wurtzite ( $B_0 = 112.3$ GPa and $B_0' = 3.5$ )				
P (GPa)	a (Å)	c (Å)	V(Å <sup>3</sup> )	Zn-S (Å)
0.0	3.79	6.14	38.14	2.312/2.316
4.0	3.74	6.06	36.64	2.282/2.286
8.0	3.70	6.00	35.47	2.258/2.261
12.0	3.66	5.95	34.47	2.236/2.239
16.0	3.63	5.90	33.60	2.218/2.219
20.0	3.60	5.86	32.83	2.201/2.202
25.0	3.57	5.81	32.01	2.182/2.183
30.0	3.54	5.77	31.27	2.165/2.166
35.0	3.51	5.74	30.61	2.150/2.150
Rock salt ( $B_0 = 155.4$ GPa and $B_0' = 5.1$ )				
P (GPa)	a (Å)		V(Å <sup>3</sup> )	Zn-S (Å)
0.0	3.54		31.48	2.510
4.0	3.51		30.50	2.480
8.0	3.49		29.93	2.464
12.0	3.46		29.36	2.449
16.0	3.44		28.80	2.433
18.0	3.43		28.52	2.425

Lattice vibrations play an important role for materials modeling, and their behavior under pressure provides useful information regarding structural instabilities and phase transformations. The frequencies ( $\omega$ ) of Raman-active modes for the W structure have been calculated as well as the Grüneisen

parameters ( $\gamma=B_0\partial\ln\omega/\partial P$ ) from their pressure dependences. Figure 5 shows the shift of the corresponding frequencies for the Raman active modes of W structure as a function of pressure. Table 5 lists for W and ZB structures the calculated pressure-coefficients of all modes and their Grüneisen parameters calculated using the value of  $B_0=112.3$  GPa found in the present study. It can be seen that the W structure presents one soft mode characterized by a decrease of the vibrational frequency with pressure. This mode has  $E_2$  symmetry and is associated to an asymmetric bending between Zn-S units, suggesting that at higher pressure the hexagonal phase should undergo a transition involving a strong coupling between a zone-centre optic mode and a strain, in this case of  $E_2$  symmetry. The pressure dependence of Raman modes as well as Grüneisen parameters are according to referenced optical phonons in ZnO.<sup>72</sup>

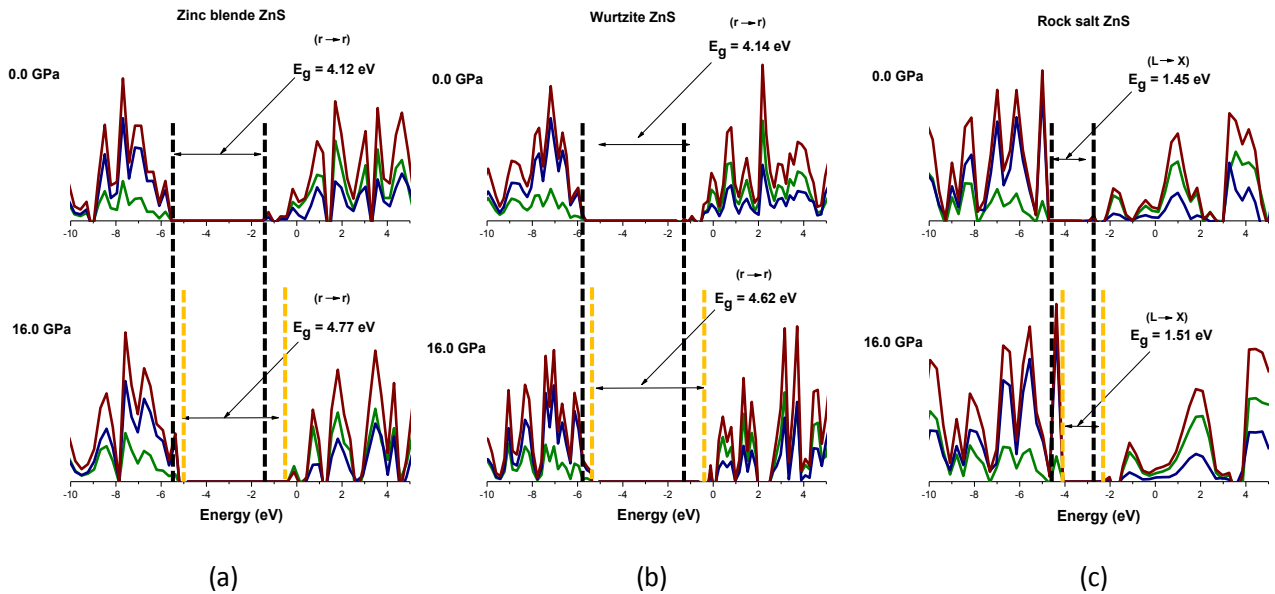
**Figure 5.** Pressure dependence of the first-order Raman modes for W structure.



**Table 5.** Calculated pressure-coefficients of Raman modes of Wurtzite and Zinc Blende, and their Grüneisen parameters.

		$\omega_0$	$\delta\omega/\delta P$	$\gamma_0$
Wurtzite	$E_2$	81.70	-0.89	-1.23
	A	315.06	4.30	1.53
	$E_1$	328.34	4.11	1.41
	$E_2$	333.08	4.46	1.50
Zinc Blende	F	358.6	4.01	1.29

The pressure dependence of the band gap of ZnS structures is shown in the Figure 6. In particular, the band gap depends of the degree of structural and electronic disorder in the lattice. The DOS analysis of ZnS polymorphs evidence changes in the VB and CB, increasing the value of the band gap with pressure. There is a major shift in the CB of ZnS polymorphs (see Figure 6) and it is possible to monitor the band gap behavior with the increasing pressure. These structural transformations induced by pressure causes a change in the electronic structure of the ZnS polymorphs (see Figure 6). Similar findings has been reported for other systems.<sup>42,43</sup>



**Figure 6:** Projected density of states (PDOS) of ZnS polymorphs with the pressure.



According to our calculations a linear behavior is observed between the values of the band gap with increasing pressure. Our results show greater variation in the value of the band gap energies with increasing pressure to the ZB and W phases, while a small change is observed for RS phase. These results follow the same trend of the theoretical study of Gupta et al.<sup>73</sup> and Huang et al.,<sup>74</sup> but with a better description in value of the band gap of these materials and in very good agreement with the experimental results reported by Ves et al.<sup>67</sup>

#### IV. CONCLUSIONS

We have investigated the structural and electronic properties and its pressure-induced phase transitions of ZnS polymorphs by periodic DFT calculations with a long-range dispersion correction. Our results confirm that this method provides a good description of the effect of the phase transition pressure in these materials.

The main results of the present work can be summarized as follows: i) Enthalpy curves as a function of pressure show that both phases, W and ZB, present very similar enthalpy, and transition pressures of the W-to-RS and ZB-to-RS are very close to each other, at 15 GPa and 15.5 GPa, respectively. These phase transitions are accompanied by an increase of the first shell coordination number of Zn atom and the unit cell volume collapse of 13.9 and 14.3% for both ZB and W phases, respectively. ii) RS phase presents a bulk modulus approximately 25% larger than the value for ZB and W structures. This decrease of the bulk compressibility is caused by the rearrangement of the polyhedral units that takes place at the transition. W presents highly anisotropic behaviour under compression with larger compressibility of the *c*-axis compared to *a*-axis. iii) The DOS analysis of ZnS polymorphs evidences changes in the VB and CB with a linear increase of the band gap value with the pressure, having a major impact on CB. iv) ZB and W phases a higher variation in the band gap energies with increasing pressure

than for RS phase, according to other theoretical and experimental studies. v) The calculated binding energies for the Zn 3*d* band is consistent with experimental XPS data. v) from the analysis of the atomic contributions of CB and VB, we obtain that the chemical bond between sulfur and the zinc atom in the ZnS polymorphs is predominantly covalent for W and ZB phases and more ionic for RS one. vi) The characterization of Raman-active vibrational frequencies for W and ZB structures and the Grüneisen parameters show a soft mode with symmetry  $E_2$  for W phase, associated to an asymmetric bending between Zn-S units, suggesting that at higher pressure the hexagonal phase should undergo reconstructive phase transition with a strong coupling between a zone-centre optic mode and a strain, in this case of  $E_2$  symmetry.

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